# **Half-linear Euler type differential equation with periodic coefficients**

#### Ondˇrej Došlý, Brno, Czech Republic

Masaryk University, Department of Mathematics and Statistics

<span id="page-0-0"></span>Malá Morávka, 29.3.2014

#### **Contents**





- **[Half-linear Euler equation](#page-11-0)**
- **[Modified Riccati equation](#page-15-0)**
- **[Periodic perturbations](#page-23-0)**

### **Half-linear differential equations**

Differential equation with the one-dimensional *p*-Laplacian

$$
(HL) \qquad (r(t)\Phi(x'))' + c(t)\Phi(x) = 0, \quad \Phi(x) = |x|^{p-2}x, \ p > 1,
$$

*r*, *c* continuous functions,  $r(t) > 0$ , Special case  $p = 2$ 

$$
(SL) \hspace{1cm} (r(t)x')' + c(t)x = 0.
$$

<span id="page-2-0"></span>linear Sturm-Liouville differential equation

### **Why 1/2-linear equations**

Motivation

Partial differential equations with *N*-dimensional *p*-Laplacian

$$
\Delta_{\rho} u + c(x) \Phi(u) = 0, \quad \Delta_{\rho} u = \text{div}\,\left(||\nabla u||^{p-2} \nabla u\right)
$$

with spherically symmetric potential *c*, i.e.,  $c(x) = c(||x||)$ , can be reduced to (HL).

If  $\|x\| = t$  then

$$
\Delta_p u(x) = t^{1-N} (t^{N-1} \Phi(u'(t)))'
$$

• Non-Newtonian fluid theory, models in glaceology,...

Equation (HL) is the Euler-Lagrange equation of the *p*-degree functional

$$
\mathcal{F}(y; a, b) = \int_{a}^{b} \left[ r(t)|y'(t)|^{p} - c(t)|y(t)|^{p} \right] dt
$$

**•** Extension of the results for the *linear* equation (special case  $p = 2$ in (HL))

$$
(SL) \hspace{1cm} (r(t)x')' + c(t)x = 0.
$$

to (HL).

• Pioneering works: I. Bihari (1963, 1968), Á. Elbert (1979), D. Mirzov (1976).

#### **Differences between linear and half-linear**

Differences between linear and half-linear.

*h*

- The solution space is only homogeneous, but not additive  $\implies$ half-linear equations.
- $\triangleright$  No half-linear analogue of the linear transformation identity

$$
h [(rx')' + cx] \stackrel{x = hy}{\rightarrow} = (By')' + Cy,
$$
  

$$
R = rh^2, \quad C = h [(rh')' + ch]
$$

 $\triangleright$  No reduction of order formula (D'Alembert formula) :  $p = 2$  and  $x(t) \neq 0$  is a solution of (SL)  $\implies$ 

$$
\tilde{x}(t) = x(t) \int_{}^t \frac{ds}{r(s)x^2(s)}
$$

is also a solution of (SL).

 $\triangleright$  No problems with the existence, uniqueness and continuability of solutions, no singular solutions, in contrast to the Emden-Fowler differential equation (which has not homogeneity property!)

$$
x''+c(t)|x|^{p-2}x=0,
$$

where no uniqueness and no continuability is guaranteed (singular solution of the first and second kind).

### **Half-linear trigonometric functions**

Half-linear trigonometric functions.

• Denote by  $S = S(t)$  the solution of

$$
(\Phi(x'))' + (p-1)\Phi(x) = 0, \quad x(0) = 0, \ x'(0) = 1.
$$

Multiplying by S' and using the initial condition

$$
|S(t)|^p + |S'(t)|^p = 1 \quad \implies \quad S' = \sqrt[p]{1 - S^p}
$$

for  $t > 0$  small. Further denote

$$
\pi_{\boldsymbol{\rho}} := 2\int_0^1 (1-s^{\boldsymbol{\rho}})^{-\frac{1}{\boldsymbol{\rho}}} \, d{\bf s} = \frac{2}{\boldsymbol{\rho}} B(1/\boldsymbol{\rho}, 1/q) = \frac{2\pi}{\boldsymbol{\rho}\sin\frac{\pi}{\boldsymbol{\rho}}}
$$

and let

$$
\sin_p t := "2\pi_p \text{ periodic odd continuation of } S(t)",
$$
  

$$
\cos_p t := (\sin_p t)'
$$

#### **Half-linear Prüfer transformation**

#### Half-linear Prüfer transformation

$$
x(t) = \rho(t) \sin_p \varphi(t), \quad r^{q-1}(t) x'(t) = \rho(t) \cos_p \varphi(t),
$$

$$
\varphi' = r^{1-q}(t) |\cos_p \varphi|^p + \frac{c(t)}{p-1} |\sin_p \varphi|^p,
$$
  

$$
\rho' = \Phi(\sin_p \varphi(t)) \cos_p \varphi(t) \left[ r^{1-q}(t) - \frac{c(t)}{p-1} \right] \rho.
$$

• The right-hand side of the equation for  $\varphi$  is Lipschitzian with respect to  $\varphi$  (and does not contain  $\rho$ )  $\implies$  existence, uniqueness, and continuability for  $\varphi$ ,  $\rho \implies$  existence, uniqueness and continuability for (HL).

#### **Oscillation theory**

- Linear Sturmian separation and comparison theory extends (almost) verbatim to (HL).
- $\triangleright$  Half-linear Riccati equation for  $w = r\Phi(x'/x)$ :

$$
(RE) \t w' + c(t) + (p-1)r^{1-q}(t)|w|^q = 0, \quad \frac{1}{p} + \frac{1}{q} = 1
$$

 $\triangleright$  Associated energy functional

<span id="page-9-0"></span>
$$
(F) \hspace{1cm} \mathcal{F}(y) = \int_{a}^{b} [r(t)|y'|^{p} - c(t)|y|^{p}] dt.
$$

### **Roundabout theorem**

The following statements are equivalent:

- (HL) is disconjugate on [*a*, *b*], i.e., every nontrivial solution has at most one zero in [*a*, *b*].
- There exists a solution of the Riccati equation

$$
w' + c(t) + (p-1)r^{1-q}(t)|w|^q = 0
$$

defined on the whole interval [*a*, *b*].

• The energy functional

$$
\mathcal{F}(y; a, b) = \int_a^b [r(t)|y'(t)|^p - c(t)|y(t)|^p] dt
$$

is positive for every nontrivial  $y \in W^{1,p}(a,b)$  with  $y(a) = 0 = y(b)$ .

#### **Half-linear Euler differential equation**

• The half-linear Euler equation

$$
(\Phi(x'))' + \frac{\gamma}{t^p}\Phi(x) = 0
$$

is oscillatory if  $\gamma > \gamma_{\bm p} := \left( \frac{p-1}{\bm p} \right)$  $\left(\frac{-1}{\rho}\right)^{\rho}$  and nonoscillatory in the opposite case.

- If  $\gamma = \gamma_{\bm p}$ , then  $x(t) = t^{\frac{\bm p-1}{\bm p}}$  is a solution
- The potential *t<sup>-p</sup>* is a border line between oscillation and nonoscillation, Kneser type (non)oscillation criteria for the equation

<span id="page-11-0"></span>
$$
\left(\Phi(x')\right)' + c(t)\Phi(x) = 0
$$
  
 
$$
\liminf_{t \to \infty} t^p c(t) > \gamma_p, \quad \limsup_{t \to \infty} t^p c(t) < \gamma_p.
$$

## **Half-linear conditional oscillation**

#### • The equation

$$
(HL) \qquad \qquad (r(t)\Phi(x'))' + c(t)\Phi(x) = 0
$$

with positive *c* is conditionally oscillatory if there exists  $\lambda_0 > 0$ , the so-called oscillation constant, such that the equation

$$
(r(t)\Phi(x'))' + \lambda c(t)\Phi(x) = 0
$$

is oscillatory for  $\lambda > \lambda_0$  and nonoscillatory for  $\lambda < \lambda_0$ . If  $\int^{\infty}$   $r^{1-q}(t)$   $dt = \infty$   $(\frac{1}{\rho} + \frac{1}{q} = 1)$ , then the equation

$$
(r(t)\Phi(x'))' + \frac{1}{r^{q-1}(t)(\int^t r^{1-q}(s) \, ds)^p}\Phi(x) = 0
$$

is conditionally oscillatory with the oscillation constant  $\lambda_0 = \gamma_p$ .

### **Limiting case**

What happens when

$$
\lim_{t\to\infty} t^{\rho}c(t)=\gamma_{\rho}:=\left(\frac{\rho-1}{\rho}\right)^{\rho},\,
$$

i.e., we have

$$
c(t) = \frac{\gamma_p}{t^p} + d(t) \quad \implies \quad (\Phi(x'))' + \left[\frac{\gamma_p}{t^p} + d(t)\right] \Phi(x) = 0
$$

with a "small" function *d*. This motivates the investigation of various perturbations of the "critical" half-linear Euler differential equation

(EE) 
$$
(\Phi(x'))' + \frac{\gamma_p}{t^p} \Phi(x) = 0.
$$

**•** Transformation approach in the linear case

$$
(*) \quad x'' + \left[\frac{1}{4t^2} + d(t)\right]x = 0 \quad \left|x = \sqrt{t}y\right| \quad (ty')' + td(t)y = 0.
$$

• The change of independent variable  $s = \log t$ ,  $y(s) = x(t)$ , in the last equation gives

$$
\frac{d^2}{ds^2}y(s)+t^2d(t)y(s)=0
$$

and from Euler equation we know that the limiting case is

$$
t^2d(t)=\frac{1}{4s^2}\quad \implies \quad d(t)=\frac{1}{4t^2\log^2 t}.
$$

o If

$$
\liminf_{t\to\infty}t^2\log^2 t\,d(t)>\frac{1}{4}
$$

then (\*) is oscillatory, if lim sup $_{t\to\infty}$   $t^2$  log $^2$   $t$   $d(t) < \frac{1}{4}$  $\frac{1}{4}$ , then (\*) is nonoscillatory.

**[Introduction](#page-2-0) [Oscillation theory](#page-9-0) [Half-linear Euler equation](#page-11-0) [Modified Riccati equation](#page-15-0) [Periodic perturbations](#page-23-0)**

# **Modified Riccati equation**

In the linear case, the transformation  $x = h(t)y$  transforms the equation

$$
(r(t)x')' + c(t)x = 0
$$

into the equation

<span id="page-15-0"></span>
$$
(R(t)y')' + C(t)y = 0
$$

with  $R = rh^2$ ,  $C = h[(rh')' + ch]$ . In terms of the associated Riccati equations

$$
w=\frac{rx'}{x}=\frac{r(h'y+hy')}{hy}=\frac{rh'}{h}+\frac{1}{h^2}\frac{rh^2y'}{y},
$$

hence

$$
v=h^2(w-w_h), v=\frac{rh^2y'}{y}, w_h=\frac{rh'}{h}.
$$

*w*

This motivates the tranformation of the Riccati equation associated with (HL)

$$
w' + c(t) + (p-1)r^{1-q}(t)|w|^q = 0, \quad \frac{1}{p} + \frac{1}{q} = 1
$$

of the form

$$
v(t) = h^{p}(t)(w(t) - wh(t)), \quad wh = \frac{r\Phi(h')}{h}.
$$

By a direct computation *v* satisfies the Modified Riccati equation:

$$
(MRE) \tV' + h^{p}(t)C(t) + (p - 1)r^{1-q}(t)h^{-q}(t)H(t, v) = 0,
$$
  
\n
$$
H(t, v) := |v + G(t)|^{q} - qv\Phi^{-1}(G(t)) - |G(t)|^{q},
$$
  
\n
$$
G(t) := r(t)h(t)\Phi(h'(t)), \quad C = h[(r\Phi(h'))' + c\Phi(h)]
$$

Special cases

 $p = 2 \implies$ 

$$
H(t, v) = (v + G)^2 - 2Gv - G^2 = v^2,
$$

hence modified RE is the equation corresponding to the transformed equation

$$
(rh^2y')' + h[(rh')' + ch]y = 0.
$$
\n
$$
\triangleright
$$
 If  $r(t) = 1$ ,  $h(t) = t^{\frac{p-1}{p}}$ , then  $G(t) \equiv \left(\frac{p-1}{p}\right)^{p-1} =: \Gamma_p$  and\n
$$
H(v, G) = |v + \Gamma_p|^q - v + \gamma_p
$$
\nand\n
$$
G(t) = t^{p-1} \left(\frac{g(t)}{p}\right)^{p}
$$

$$
C(t)=t^{p-1}\left(c(t)-\frac{\gamma_p}{t^p}\right).
$$

Quadratization of the function H: Suppose that  $h'(t) \neq 0$ ,

$$
H(t, v) = |G|^q \left\{ \left| \frac{v}{G} + 1 \right|^q - q\frac{v}{G} - 1 \right\}
$$
  
 
$$
\sim \frac{q(q-1)}{2} |G|^q \left( \frac{v}{G} \right)^2 = \frac{q(q-1)}{2} |G|^{q-2} v^2 \text{ as } v \to 0.
$$

We have

$$
(p-1)r^{1-q}h^{-q}H(t,v)\sim \frac{q}{2}\frac{v^2}{R}, \quad R:=rh^2|h'|^{p-2}.
$$

hence we obtain the Approximate Riccati equation

$$
(ARE) \t v' + C(t) + \frac{q}{2} \frac{v^2}{R(t)} = 0.
$$

(ARE) is the Riccati equation associated with the linear Sturm-Liouville equation

$$
(R(t)y')' + \frac{q}{2}C(t)y = 0
$$

Half-linear Euler type differential equation with periodic coefficients

.

# **Application of MRE**

Elbert and Schneider, 2000. The half-linear Riemann-Weber equation

$$
(\Phi(x'))' + \left(\frac{\gamma_p}{t^p} + \frac{\mu}{t^p \log^2 t}\right) \Phi(x) = 0, \quad \gamma_p = \left(\frac{p-1}{p}\right)^p,
$$

is oscillatory iff

$$
\mu > \mu_p := \frac{1}{2} \left( \frac{p-1}{p} \right)^{p-1}
$$

Modified Riccati equation is (with  $h(t) = t^{\frac{\rho-1}{\rho}}$ )

$$
v' + \frac{\mu}{t \log^2 t} + (\rho - 1)t^{-1} [ |v + \Gamma_{\rho}|^q - v + \gamma_{\rho}] = 0
$$

and the Approximate Riccati equation is

$$
z' + \frac{\mu}{t \log^2} + \frac{1}{2t\Gamma_p}z^2 = 0
$$

The change of independent variable  $s = \log t$  gives

$$
z' + \frac{\mu}{s^2} + \frac{1}{2\Gamma_p}z^2 = 0.
$$

The associated second order differential equation is

$$
\left(2\Gamma_\rho y'\right)' + \frac{\mu}{s^2}y = 0
$$

which the same as

$$
y''+\frac{\mu}{2\Gamma_p}\frac{1}{s^2}y=0
$$

and the last equation is nonoscillatory iff

$$
\frac{\mu}{2\Gamma_{\rho}} \leq \frac{1}{4} \quad \iff \quad \mu \leq \frac{1}{2}\Gamma_{\rho} = \frac{1}{2}\left(\frac{\rho - 1}{\rho}\right)^{\rho - 1} =: \mu_{\rho}.
$$

More general perurbation of the critical Euler equation

$$
\left(\Phi(x')\right)' + \frac{1}{t^p} \left[\gamma_p + \frac{\mu_p}{\log^2 t} + \frac{\mu}{\log^2 t \log^2(\log t)}\right] \Phi(x) = 0
$$

is using the same method nonoscillatory iff  $\mu \leq \mu_p$ . We consider a more general perturbation of critical Euler equation (EE):

$$
\text{(EP)} \quad \left[ \left( 1 + \sum_{j=1}^{n} \frac{\alpha_j}{\text{Log}_j^2(t)} \right) \Phi(x') \right]' + \left[ \frac{\gamma_p}{t^p} + \sum_{j=1}^{n} \frac{\beta_j}{t^p \text{Log}_j^2(t)} \right] \Phi(x) = 0
$$

with

$$
Log_k(t) = \prod_{j=1}^k log_k(t), \quad log_k(t) = log_{k-1}(log t), \ log_1(t) = log t.
$$

# **(Non)oscillation of (EP)**

Recall that

$$
\gamma_p = \left(\frac{p-1}{p}\right)^p, \quad \mu_p = \frac{1}{2}\left(\frac{p-1}{p}\right)^p.
$$

O. D. H. Funková (2012):

- **If**  $\beta_1 + \gamma_0 \alpha_1 < \mu_0$  then (EP) is nonoscillatory and if  $\beta_1 + \gamma_0 \alpha_1 > \mu_0$ then it is oscillatory.
- Let  $\beta_1 + \gamma_0 \alpha = \mu_0$ , if  $\beta_2 + \gamma_0 \alpha_2 < \mu_0$  then (EP) is nonoscillaotry and if  $\beta_2 + \gamma_0 \alpha_2 > \mu_0$  then it is oscillatory.

Let

$$
\beta_k + \gamma_p \alpha_k = \mu_p, \quad k = 1, \ldots, n-1.
$$

. . .

Then (EP) is nonoscillatory if and only if

$$
\beta_n + \gamma_p \alpha_n \leq \mu_p.
$$

### **Periodic perturbations**

We consider the equation

$$
(HLP) \qquad \qquad (r(t)\Phi(x'))' + \frac{\lambda c(t)}{t^p}\Phi(x) = 0
$$

with positive α-periodic functions *r*, *c*

- **If**  $r(t) \equiv 1$ ,  $c(t) \equiv 1$ , then (HLP) reduces to the half-linear Euler equation and its oscillation constant is  $\lambda_\mathbf{0} = \gamma_{\boldsymbol{\rho}} = \left( \frac{\boldsymbol{\rho}-\mathbf{1}}{\boldsymbol{\rho}} \right)$  $\left(\frac{-1}{p}\right)^p$ .
- **Theorem** (P. Hasil, 2009). Let

$$
\overline{r}=\frac{1}{\alpha}\int_0^\alpha\frac{ds}{r^{q-1}(s)},\quad \overline{c}=\frac{1}{\alpha}\int_0^\alpha c(s)\,ds.
$$

Then (HLP) is oscillatory if

<span id="page-23-0"></span>
$$
\lambda > \gamma_{r,c} := \frac{\gamma_p}{(\overline{r})^{p-1}\overline{c}}.
$$

and nonoscillatory if  $\lambda < \gamma_{r,c}$ .

# **The limiting case**  $\lambda = \gamma_{r,c}$

Consider the equation

(RW) 
$$
(r(t)\Phi(x'))' + \left[\frac{\gamma_{r,c}c(t)}{t^p} + \frac{\mu d(t)}{t^p \log^2 t}\right]\Phi(x) = 0.
$$

**Theorem** (O. D., P. Hasil, 2011). Let *r*, *c* be as before and

$$
\overline{d}=\frac{1}{\alpha}\int_0^\alpha d(s)\,ds.
$$

Then (RW) is oscillatory if

$$
\mu > \mu_{r,d} := \frac{\mu_p}{(\overline{r})^{p-1}\overline{d}}
$$

and nonoscillatory if  $\mu < \mu_{r,d}$ . In particular, equation (HLP) is nonoscillatory also in the limiting case  $\lambda = \gamma_{r,c}$ .

#### Consider the equation (EPP)

$$
\left[\left(r(t)+\sum_{j=1}^n\frac{\alpha_j(t)}{\text{Log}_j^2(t)}\right)^{1-p}\Phi(x')\right]'+\left[\frac{c(t)}{t^p}+\sum_{j=1}^n\frac{\beta_j(t)}{t^p\text{Log}_j^2(t)}\right]\Phi(x)=0.
$$

with  ${\mathcal T}$  periodic functions  $\alpha_j, \beta_j, j=1,\ldots,n.$  Denote by  $\bar r, \bar c, \, \bar \alpha_j, \bar \beta_j$  $j = 1, \ldots, n$ , their mean values over the period  $T$ , i.e.,

$$
\bar{r} = \frac{1}{T} \int_0^T r(t) dt, \quad \bar{c} = \frac{1}{T} \int_0^T c(t) dt,
$$
  

$$
\bar{\alpha}_j = \frac{1}{T} \int_0^T \alpha_j(t) dt, \quad \bar{\beta}_j = \frac{1}{T} \int_0^T \beta_j(t) dt.
$$

**Theorem.** (O. D. H.Funková, 2013). Let  $\bar{c}\bar{r}^{p-1} = \gamma_p$ . If there exists  $k \in \{1, \ldots, n\}$  such that

$$
\bar{\beta}_j \bar{r}^{p-1} + (p-1)\gamma_p \bar{\alpha}_j \bar{r}^{-1} = \mu_p, \quad j = 1, \ldots, k-1,
$$

and  $\bar{\beta}_k \bar{r}^{p-1} + (p-1) \gamma_p \bar{\alpha}_k \bar{r}^{-1} \neq \mu_p,$  then (EPP) is oscillatory if

$$
\bar{\beta}_k \bar{r}^{p-1} + (p-1)\gamma_p \bar{\alpha}_k \bar{r}^{-1} > \mu_p
$$

and nonoscillatory if

$$
\bar{\beta}_k \bar{r}^{p-1} + (p-1)\gamma_p \bar{\alpha}_k \bar{r}^{-1} < \mu_p.
$$

- O. DOŠLÝ, H. FUNKOVÁ, *Perturbations of half-linear Euler differential equation and transformations of modified Riccati equation.* Abstr. Appl. Anal. 2012, Art. ID 738472, 19 pp.
- E. O. DOŠLÝ, H. FUNKOVÁ, *Euler type half-linear differential equations with periodic coefficients.* Abstr. Appl. Anal. 2013, Art. ID 714263, 6 pp.
- 譶 O. DOŠLÝ, P.HASIL, *Critical oscillation constant for half-linear differential equations with periodic coefficients*, Annal. Mat. Pura Appl. **190** (2011), no. 3, 395–408.
- O. DOŠLÝ, P. ŘEHÁK Half-linear Differential Equations, F North-Holland Mathematics Studies, 202. Elsevier Science B.V., Amsterdam, 2005. xiv+517 pp.
- Á. ELBERT, A. SCHNEIDER , *Perturbations of the half-linear Euler* F *differential equation*. Results Math. **37** (2000), no. 1-2, 56–83.
- Ħ P. HASIL, *Conditional oscillation of half-linear differential equations with periodic coefficients*, Arch. Math. (Brno) **4**4 (2008), 119–131.
- F P. Hasil, M. Veselý, *Oscillation of half-linear differential equations with asymptotically almost periodic coefficients*. Adv. Difference Equ. 2013, 2013:122, 15 pp.
- P. HASIL, *Conditional oscillation of half-linear differential equations* F. *with periodic coefficients*, Arch. Math. (Brno) **4**4 (2008), 119–131.
- R H. KRÜGER, G. TESCHL, *Effective Prüfer angles and relative oscillation criteria*, J. Differential Equations **245** (2009), 3823–3848.
- <span id="page-28-0"></span>K. M. SCHMIDT, *Critical coupling constant and eigenvalue* 暈 *asymptotics of perturbed periodic Sturm-Liouville operators*, Commun Math. Phys. **211** (2000), 465–485.